

Rate Constants of HO₂ + NO Covering Atmospheric Conditions. 1. HO₂ Formed by OH + H₂O₂

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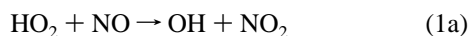
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Rate constants of the gas-phase reactions HO₂ + NO → products (1), HO₂ + NO → OH + NO₂ (1a), and OH + NO → products (2) were determined at room temperature and total pressures of 10, 50, and 100 kPa of N₂. OH radicals were produced by pulsed 248 nm photolysis of H₂O₂ and monitored by time-resolved CW UV-laser long-path absorption. Biexponential OH decay curves, observed in the presence of H₂O₂ and NO, were analyzed to obtain rate constants for the above reactions. Those of reactions 1 and 1a were found to be independent of pressure with averaged values of $(9.7 \pm 1.5) \times 10^{-12}$ and $(9.6 \pm 1.5) \times 10^{-12}$ cm³ s⁻¹, respectively. This indicates an OH yield very close to unity (≥ 0.95) for the title reaction under the experimental conditions of this work. Error limits are estimated considering statistical and possible systematic errors. The pressure dependence found for reaction 2 is in good agreement with literature data. Rate constants of $(1.3 \pm 0.4) \times 10^{-12}$, $(4.7 \pm 0.8) \times 10^{-12}$, and $(7.4 \pm 1.3) \times 10^{-12}$ cm³ s⁻¹ have been determined for total pressures of 10, 50, and 100 kPa of N₂, respectively. Addition of H₂O at mixing ratios of 1.7% had no influence on the investigated reactions. Moreover, a quantum yield ≤ 0.05 is estimated for the formation of H + HO₂ in the 248 nm photolysis of H₂O₂. The rate constants of the HO₂ + NO reaction are slightly larger (15%) than the most recent recommendations.

Introduction

The HO₂ radical is an important species in atmospheric and combustion chemistry. In the troposphere it mainly reacts with NO forming OH and NO₂.¹



The photolysis of NO₂ then leads to the formation of ozone.¹ The HO₂ + NO reaction has been investigated in a large number of studies before, and there is now good agreement on the rate constant at low pressure conditions.^{2–11} On the other hand, studies of reaction 1 have not been performed at total pressures exceeding ≈ 2 kPa until recently.¹¹ The slightly negative temperature coefficient of the reaction^{2–4,11} indicates the possibility of a transient complex formation^{10,12} which might cause a pressure dependence of the rate constant (like, for example, in the HO₂ self-reaction¹³). Thus, a pressure-dependent study of the HO₂ + NO reaction covering the tropospheric pressure range was needed.

Two other reactions might compete with (1a) and have been discussed in the literature.^{2,10}



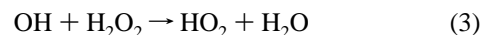
Up to now, no indication for one of these reactions has been reported. The product of reaction 1b, peroxyxynitrous acid (HOONO), has been identified in argon matrices as a product of nitric acid photolysis¹⁴ but could not be detected as a product of the OH + NO₂ reaction in the gas phase.¹⁵ Alkyl peroxy radicals are known to react with NO forming an adduct with subsequent isomerization to the corresponding nitrates.¹⁶ However, this process is only of importance for species containing

more than three C atoms. Thus, although the energy released in the HO₂ + NO reaction is roughly a factor of 2 smaller,¹⁷ a stabilization of HOONO, possibly followed by an isomerization into nitric acid, appears to be unlikely at atmospheric pressure conditions. Nevertheless, it could not be excluded totally and may be of importance at very high pressures.

In this work we have performed measurements at total pressures of 10, 50, and 100 kPa of N₂. The experimental method employed here allows us to determine the rate constants of reactions 1 and 1a, as well as that of the reaction of OH with NO, forming HONO.^{18,19}



H₂O₂ is photolyzed in a 248 nm excimer laser pulse to produce ground state OH radicals. The reaction of OH with H₂O₂ forms HO₂



from which OH is regenerated in the presence of NO by reaction 1a. This results in biexponential OH decay curves that were analyzed to obtain the rate constants of the involved reactions.

Experimental Section

OH is monitored by time-resolved detection by CW UV-laser long-path absorption on the Q₁(2) line of the OH(A←X) transition around 308 nm (path length 100 m, CW laser power ≈ 200 μW). This experimental method has been described before.^{20–22} H₂O₂ concentrations varied between 1.5 and 4.5 Pa in a large excess of N₂. With the available excimer laser pulse energy (≈ 300 mJ in a beam diameter of 15 cm) this results in typical starting concentrations of about 5×10^{10} cm⁻³ of OH in a gas volume of approximately 20 L. This concentration is sufficient to obtain reasonable decay curves after averaging

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over 10–20 laser shots while it is small enough to ensure pseudo-first-order conditions.

The measurements were carried out at room temperature in slowly flowing gas mixtures at total pressures of 10, 50, and 100 kPa of N₂ containing 5, 0.4, and 0.1% of O₂, respectively. O₂ was added in order to convert H atoms, possibly formed in the H₂O₂ photolysis, into HO₂. The gases had minimum purities of 99.999% (N₂), 99.995% (O₂), and 99.5% (NO) stated by the manufacturer (Messer Griesheim). The combination of mass-flow controllers (FC 260, Tylan) with capacitance manometers (Baratron, MKS) allows us to adjust concentrations and total pressure. NO pressures were varied between 0.2 and 0.6 Pa. Diluted gas mixtures containing 300 or 2000 ppm of NO in N₂ or Ar were used. The concentrations have stated uncertainties of 2% (Messer Griesheim). Before entering the cell these mixtures were fed through a glass tube filled with solid iron(II) sulfate (FeSO₄·5H₂O) in order to remove NO₂ traces.

Measurements at a given total pressure were made with at least two significantly different H₂O₂ concentrations in the range given above. The dosage of H₂O₂ by purging a concentrated liquid solution with inert gas, as well as its limited stability in the reaction cell have been discussed in our recent work on the OH + C₂H₂ reaction.²² H₂O₂ concentrations were calculated by measuring the OH decay rate in the absence of other reactants (τ_0^{-1}), using a rate constant of $1.75 \times 10^{-12} \text{ cm}^3 \text{ s}^{-1}$ for the OH + H₂O₂ reaction.²³ No evidence for significant OH losses due to background impurities was found. A plot of the gas flow purging the H₂O₂ solution against τ_0^{-1} gave a linear dependence leading through the origin in good approximation. We estimate that the OH loss rate coefficient in the absence of H₂O₂ is below 10 s^{-1} for the total pressures used here.

Since NO losses are recognized in the OH time behavior even after only 10 laser shots, the reaction cycle initiated by the production of OH appears to be very effective. As a consequence, decay curves were recorded by averaging over three laser shots for the time being. Due to the large volume ($\approx 20 \text{ L}$) and the limited total gas flow ($\approx 3 \text{ L/min (STP)}$) 20 min were then allowed to elapse until the next measurement was started. This procedure was repeated about 10 times to complete a measurement.

Figure 1 shows typical decay curves obtained with different NO concentrations. The recorded signal is an amplified transient difference in photocurrents, proportional to the absorbance of the gas mixture and thus to the OH concentration at low optical densities. At $t \leq 0.1 \text{ ms}$ an increase of the signal can be recognized owing to the rise time of the amplifier acting as a low-pass filter. As a consequence, the absorption signals are recorded somewhat delayed with respect to their real time behavior. To describe the time dependence of the recorded signal the convolution integral of the filter function, with time constant τ_D , and the time dependence of the OH concentration has to be calculated. In case of a biexponential OH time dependence

$$[\text{OH}](t)/[\text{OH}]_{t=0} = c_1 \exp(-t/\tau_1) + c_2 \exp(-t/\tau_2) \quad (4)$$

which will be discussed in more detail below, the convolution integral results in the following expression for the recorded signal.

$$\Delta I_{\text{rec}}(t) \propto \frac{c_1}{1 - \tau_D/\tau_1} \exp(-t/\tau_1) + \frac{c_2}{1 - \tau_D/\tau_2} \exp(-t/\tau_2) - \left(\frac{c_1}{1 - \tau_D/\tau_1} + \frac{c_2}{1 - \tau_D/\tau_2} \right) \exp(-t/\tau_D) \quad (5)$$

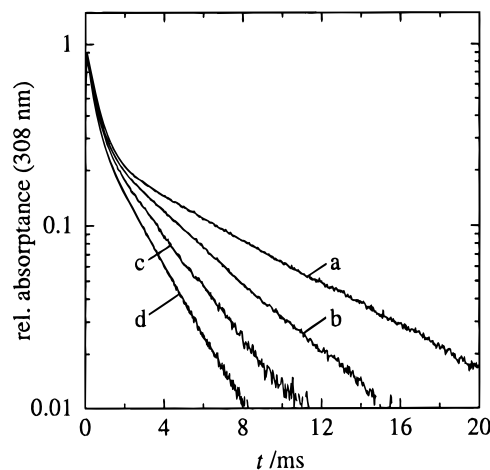
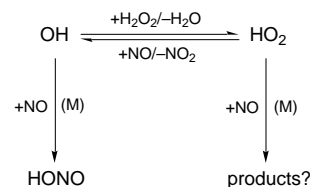


Figure 1. Typical biexponential OH decay curves in a semilogarithmic plot obtained at a total pressure of 100 kPa of N₂, an H₂O₂ concentration of $5.6 \times 10^{14} \text{ cm}^{-3}$ and different NO levels. [NO]/ 10^{14} cm^{-3} : (a) 0.60, (b) 0.85, (c) 1.09, and (d) 1.34. The displayed curves are averaged over 30 single measurements (30 laser shots) and normalized to the same extrapolated amplitudes at $t = 0$. At very short times ($\leq 0.1 \text{ ms}$) a rise of the absorption signal can be recognized, caused by the time constant of the detection system.

The time constant τ_D has been determined to be $25 \pm 5 \mu\text{s}$ in a nonlinear least squares fit to decay curves using the highest possible time resolution ($10 \mu\text{s}$) and carefully avoiding stray light from the excimer laser. For $t \geq 0.2 \text{ ms}$ the last term of eq 5 vanishes, and a biexponential time dependence with the same time constants as in eq 4 remains. Thus, the experimentally recorded decay curves were analyzed for $t \geq 0.2 \text{ ms}$ and referred to as biexponential in a nonlinear least-squares fit with amplitudes that were corrected according to eq 5, i.e., multiplied by $(1 - \tau_D/\tau_i)$. In the fitting procedure five parameters have been considered: the amplitudes and time constants of the two exponentials (eq 4), and the background signal. Corrections of the amplitude of the exponential with the smaller time constant (usually referred to as 1) are in the range of 4–8%.

Results

Evaluation of Rate Constants. The following scheme shows the reactions occurring in the presence of NO and H₂O₂ after an initial production of OH (or HO₂).



The corresponding system of differential equations for the radical concentrations

$$d[\text{OH}]/dt = -{}^1k_{\text{1,OH}}[\text{OH}] + {}^1k_{\text{r,OH}}[\text{HO}_2] \quad (6)$$

$$d[\text{HO}_2]/dt = {}^1k_{\text{r,HO}_2}[\text{OH}] - {}^1k_{\text{1,HO}_2}[\text{HO}_2] \quad (7)$$

which can be solved analytically²⁴ results in a biexponential time dependence for the OH concentration according to eq 4.

The time constants τ_1 and τ_2 depend on the formation (${}^1k_{\text{f}}$) (the index 1 at the 1k values indicates pseudo-first-order rate constants) and loss (${}^1k_{\text{l}}$) rate coefficients of OH and HO₂, while the amplitudes c_1 and c_2 also depend on starting conditions. In the special case that only OH is present at $t = 0$, which will be

shown to be the case here, the following relationships can be derived.

$${}^1k_{\text{l,HO}_2} = \frac{\tau_1^{-1} + c_1/c_2 \cdot \tau_2^{-1}}{c_1/c_2 + 1} \quad (8)$$

$${}^1k_{\text{l,OH}} = \tau_1^{-1} + \tau_2^{-1} - {}^1k_{\text{l,HO}_2} \quad (9)$$

$${}^1k_{\text{f,OH}} {}^1k_{\text{f,HO}_2} = {}^1k_{\text{l,HO}_2} {}^1k_{\text{l,OH}} - \tau_1^{-1} \tau_2^{-1} \quad (10)$$

The amplitudes c_1 and c_2 appear as a ratio. By this the proportionality factor in eq 5, containing the initial OH concentration, cancels out. On the other hand, OH and HO₂ formation rate coefficients only occur as a product. Fortunately, in the present system the HO₂ formation rate coefficient can be determined independently by measuring the OH loss rate coefficient in the absence of NO (τ_0^{-1}). In case that (i) the loss of OH is only due to reaction with H₂O₂, as confirmed in the Experimental Section, and (ii) assuming an HO₂ yield of unity for the H₂O₂ + OH reaction (3),²⁵ the HO₂ formation rate coefficient is given by

$${}^1k_{\text{f,HO}_2} = k_3[\text{H}_2\text{O}_2] = \tau_0^{-1} \quad (11)$$

Thus, the OH formation rate coefficient ${}^1k_{\text{f,OH}}$ can be calculated from eqs 8–10.

$${}^1k_{\text{f,OH}} = ({}^1k_{\text{l,HO}_2} {}^1k_{\text{l,OH}} - \tau_1^{-1} \tau_2^{-1}) \tau_0 \quad (12)$$

Table 1 summarizes the biexponential curve parameters, obtained for different NO concentrations and the independently measured values of τ_0^{-1} .

HO₂ + NO Reaction. Figure 2 shows HO₂ loss and OH formation rate coefficients as a function of NO concentration. Different H₂O₂ pressures do not influence these coefficients, as expected from the reaction model. Within the scatter of the data there is neither a pressure dependence nor a difference in OH formation and HO₂ loss rate coefficients noticeable. This is confirmed in Table 2 showing the results from linear regressions of the data to obtain the rate constants k_1 and k_{1a} for the different total pressures according to the following equations.

$${}^1k_{\text{l,HO}_2} = k_1[\text{NO}] + c_{\text{l,HO}_2} \quad (13)$$

$${}^1k_{\text{f,OH}} = k_{1a}[\text{NO}] + c_{\text{f,OH}} \quad (14)$$

No indication for an HO₂ loss or OH formation not due to the reaction with NO was found. The fitted intercepts $c_{\text{l,HO}_2}$ and $c_{\text{f,OH}}$ for the coefficients were found to be slightly negative ($\approx -(30 \pm 40)\text{s}^{-1}$), but zero within their 2σ error limits. This may be taken as a hint at a small loss of NO caused by the effective chain reaction, even while proceeding with very few laser shots, as described in the Experimental Section. However, the influence of this possible loss is smaller than the 2σ error limits stated in Table 2.

The OH yields of the HO₂ + NO reaction, i.e. the fractions of the rate constants k_{1a} and k_1 , are also given in Table 2. They were obtained by linear regressions from plots of the rate coefficients from eqs 13 and 14 against each other. The OH yields are unity in all cases with 2σ errors limits $\leq 5\%$.

Filled circles in Figure 2 indicate measurements made in the presence of water vapor with a mixing ratio of 0.017. Since

TABLE 1: NO Concentrations, Fitted Decay Curve Parameters, τ_1^{-1} , τ_2^{-1} , and c_1/c_2 (Corrected According to Eq 5), and OH Decay Rates, τ_0^{-1} (Measured in the Absence of NO)

[NO]/10 ¹⁴ cm ⁻³	$\tau_1^{-1}/10^3$ s ⁻¹	$\tau_2^{-1}/10^2$ s ⁻¹	c_1/c_2	$\tau_0^{-1}/10^3$ s ⁻¹
<i>p</i> = 9.9 kPa				
0.55	1.56	0.30	2.23	1.01
0.62	1.67	0.35	2.05	1.03
0.73	1.34	0.52	1.05	0.61
0.74	1.72	0.45	1.65	1.00
0.90	1.95	0.57	1.44	1.03
1.02	2.02	0.68	1.24	1.00
1.12	1.69	0.93	0.69	0.61
1.17	2.20	0.86	1.08	1.02
1.28	2.31	0.97	0.97	1.01
1.42	2.42	1.13	0.86	0.99
0.70 ^a	1.41	0.49	1.23	0.71
1.03 ^a	1.74	0.89	0.82	0.71
<i>p</i> = 50.0 kPa				
0.35	1.14	0.46	3.16	0.74
0.56	1.81	0.77	3.00	1.15
0.57	1.68	0.82	2.73	1.01
0.59	1.44	1.06	2.02	0.73
0.78	1.64	1.56	1.64	0.74
0.81	2.25	1.28	2.30	1.27
0.87	2.45	1.28	2.61	1.45
1.00	2.32	1.87	1.89	1.15
1.05	2.31	2.11	1.65	1.08
1.20	2.89	2.14	2.00	1.44
1.27	2.78	2.52	1.71	1.29
1.40	2.97	2.85	1.63	1.30
0.52 ^a	1.83	0.64	3.30	1.21
0.85 ^a	2.24	1.69	2.32	1.21
<i>p</i> = 99.8 kPa				
0.59	2.70	0.94	5.08	1.89
0.60	1.83	1.36	3.02	0.99
0.71	2.46	1.28	3.54	1.44
0.77	2.99	1.38	4.13	1.88
0.85	2.17	2.23	2.45	0.99
0.96	2.83	2.11	2.90	1.47
1.08	3.39	2.40	3.08	1.87
1.09	2.42	3.27	2.07	0.98
1.21	3.25	3.32	2.56	1.53
1.30	3.71	3.10	2.74	1.87
1.34	2.79	4.39	1.84	0.98
1.14 ^a	3.00	3.01	2.57	1.40
1.41 ^a	3.37	4.10	2.24	1.41

^a Measurements in the presence of 1.7% of water vapor, not considered in the evaluation of the rate constants in Table 2.

no effect was observed, no further measurements were performed in the presence of water.

OH + NO Reaction. As expected from the reaction model, the OH loss rate coefficient was found to depend on total pressure and NO and H₂O₂ concentration.

$${}^1k_{\text{l,OH}} = k_2[\text{NO}] + k_3[\text{H}_2\text{O}_2] + c_{\text{l,OH}} \quad (15)$$

In order to determine the rate constant k_2 , measurements with different H₂O₂ concentrations were combined by subtracting the τ_0^{-1} values measured in the absence of NO (eq 11). The result is shown in Figure 3 for the different total pressures used here. The straight lines correspond to the rate constants given in Table 2. The intercepts $c_{\text{l,OH}}$ were found to be zero in good approximation ($\approx -(10 \pm 40)\text{s}^{-1}$). This is not in contradiction to the suggestion that a loss of NO is responsible for the slightly negative intercepts found for the HO₂ loss rate coefficients: NO₂ is formed when NO and H₂O₂ are consumed and NO₂ is reactive with respect to OH. However, the amount of NO₂ formed this way is not expected to have a significant influence on the present reaction system.

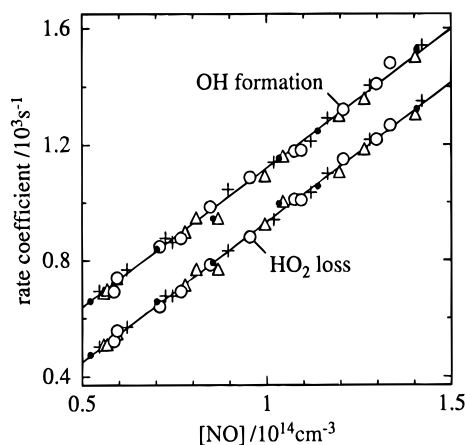


Figure 2. OH formation rate coefficient ($k_{f,\text{OH}}$) and HO₂ loss rate coefficient (k_{l,HO_2}) dependences on the NO concentration for total pressures of approximately 100 (○), 50 (△), and 10 (×) kPa total pressure of N₂. The straight lines indicate averaged rate constants from Table 2. OH formation rate coefficients are shifted by 200 s⁻¹ for clarity. Filled circles indicate measurements at water vapor mixing ratios of 0.017 at different total pressures (to assign total pressures refer to Figure 3 or Table 1).

TABLE 2: Rate Constants^a (k_1 , k_{1a}) and OH Yield^b (ϕ_{OH}) of the Reaction HO₂ + NO, and Rate Constant^a of the Reaction OH + NO (k_2) for Different Total Pressures of N₂, $T = 297 \pm 2$ K

$p(\text{N}_2)/\text{kPa}$	k_1	k_{1a}	ϕ_{OH}	k_2
9.9	9.7 ± 0.3	9.6 ± 0.4	0.99 ± 0.03	1.3 ± 0.3
50.0	9.5 ± 0.4	9.5 ± 0.3	1.00 ± 0.03	4.7 ± 0.3
99.8	9.8 ± 0.4	9.7 ± 0.6	1.00 ± 0.05	7.4 ± 0.6

^a Error limits are 2σ (statistical). Units are 10^{-12} cm³ s⁻¹, an additional relative error of 10% is estimated from uncertainties of gas-flow rates, pressure determination and further possible systematic errors.

^b Slopes taken from a correlation of the rate coefficients in Figure 2.

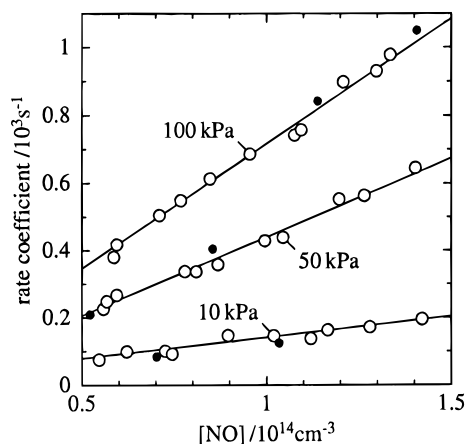


Figure 3. OH loss rate coefficients ($k_{l,\text{OH}}$) as a function of NO concentration for different total pressures of N₂. The data are normalized to $[\text{H}_2\text{O}_2] = 0$ by subtracting the OH loss rate coefficient in the absence of NO, measured in any case before and after addition of NO. The straight lines indicate the rate constants k_2 in Table 2. Filled circles indicate measurements at water vapor mixing ratios of 0.017.

The measured k_2 values are in good agreement with recent literature data.^{19,26}

H + HO₂ Formation in H₂O₂ Photolysis. Recently, a quantum yield for OH formation significantly below 2 ($\phi_{\text{OH}} = 1.58 \pm 0.23$) has been reported by Schiffman et al.²³ for the 248 nm photolysis of H₂O₂. If H + HO₂ are photolysis products formed together with OH they are expected to influence the present system. Although a significant formation of H atoms

has already been excluded in the literature,^{27,28} we will show that also the data of this work are inconsistent with this possibility. As mentioned in the Experimental Section, small amounts of O₂ were added to provide a rapid conversion of possibly formed H atoms into HO₂ ($\tau(\text{H}) \approx 30\text{--}60$ μs).²⁶ Considering the different starting conditions, eq 8 has to be modified.

$${}^1k_{l,\text{HO}_2} = \frac{1}{2f} + \frac{(\tau_1^{-1} + \tau_2^{-1})}{2} - \left[\frac{\{f(\tau_1^{-1} - \tau_2^{-1}) + 1\}^2}{4f^2} - \frac{(\tau_1^{-1} - \tau_2^{-1})}{f(c_1/c_2 + 1)} \right]^{1/2} \quad (16)$$

The quantity f depends on the ratio of (H + HO₂) and OH formed in the photolysis and the independently measured value of τ_0 , which again is taken as the reciprocal of the HO₂ formation rate coefficient.

$$f = \frac{[\text{H} + \text{HO}_2]_{t=0}}{[\text{OH}]_{t=0}} \tau_0 \quad (17)$$

Of course, in the limit $f \rightarrow 0$, eq 16 corresponds to eq 8.

All decay curves were analyzed a second time by taking the ratio $[\text{H} + \text{HO}_2]_{t=0}/[\text{OH}]_{t=0} = 0.25$ which is the expected maximum value, corresponding to an OH quantum yield of 1.6.²³ Virtually no difference in the quality of the fits was found. The NO dependences of the different coefficients were linear as in Figures 2 and 3. However, the slopes and thus the rate constants were different. The measurements at a total pressure of 100 kPa gave no indication on which attempt is correct. Rate constants ($\pm 2\sigma$) $k_1 = (7.8 \pm 0.4) \times 10^{-12}$ cm³ s⁻¹, $k_{1a} = (7.9 \pm 0.5) \times 10^{-12}$ cm³ s⁻¹, and $k_2 = (9.0 \pm 0.8) \times 10^{-12}$ cm³ s⁻¹ were derived, which are in reasonable agreement with literature data.^{17,26}

On the other hand, the results at lower total pressures are increasingly inconsistent. At 10 kPa k_2 would be a factor of 2.5 higher than recent literature values.^{19,26} Furthermore, k_1 and k_{1a} would exhibit slight pressure dependences with different directions. Between 10 and 100 kPa k_1 would increase by 12% and k_{1a} decrease by 20%. Thus, with decreasing pressure the OH yield of reaction 1 would increase to physically meaningless values (up to ≈ 1.4 at 10 kPa). From these results we estimate an upper limit of 0.05 for the quantum yield of H + HO₂ formation in the 248 nm photolysis of H₂O₂ in accordance with literature.^{27,28}

Discussion

In 1977 Howard²⁹ realized that the rate constants of the HO₂ + NO reaction measured before were too small. This was verified by Howard and Evenson² and other authors.³⁻¹¹ Table 3 summarizes room temperature measurements on the rate constant that have been made in the meantime. All data agree within the stated uncertainties except those of Thrush and Wilkinson⁸ and Rozenshtein et al.⁹ due to their relatively small error limits (1σ , not considering possible systematic errors). Up to now, all experiments have been performed in flow tubes, in most cases with LMR detection of HO₂. This technique is limited to total pressures below ≈ 2 kPa for reasons of pressure broadening. Very recently, Seeley et al.¹¹ have performed measurements at total pressures of up to 25 kPa using a turbulent flow technique and chemical ionization mass spectroscopy detection. These authors¹¹ found no pressure dependence of the rate constant in agreement with the results reported here. The data of the present study are the first measurements at

TABLE 3: Literature Data on the Rate Constant of the HO₂ + NO Reaction at Room Temperature

references	$k_1/10^{-12} \text{ cm}^3 \text{ s}^{-1}$
Howard and Evenson (1977) ²	$8.1 \pm 1.5^{a,b}$
Howard (1979) ³	$8.0 \pm 1.0^{a,b}$
Leu (1979) ⁴	$9.8 \pm 1.6^{a,b}$
Burrows et al. (1979) ⁵	8.2 ± 2.4
Glaschick-Schimpf et al. (1979) ⁶	11 ± 3^b
Hack et al. (1980) ⁷	$7.6 \pm 1.7^{a,b}$
Thrush and Wilkinson (1981) ⁸	$6.9 \pm 0.6^{a,b}$
Rozenstein et al. (1984) ⁹	7.0 ± 0.6
Jemi-Alade and Thrush (1990) ¹⁰	$8.5 \pm 1.3^{a,b}$
Seeley et al. (1996) ¹¹	8.0 ± 1.2
recommendations	
IUPAC (1992) ²⁶	$8.3 \pm \begin{smallmatrix} 2.2 \\ 1.7 \end{smallmatrix}$
NASA (1994) ¹⁷	$8.6 \pm \begin{smallmatrix} 1.7 \\ 1.4 \end{smallmatrix}$
this work	9.7 ± 1.5^c

^a Considered in IUPAC recommendation. ^b Considered in NASA recommendation. ^c Error limit estimated considering statistical (Table 2) and possible systematic errors (10%).

pressures of up to 100 kPa covering the entire tropospheric pressure range.

The current IUPAC²⁶ and NASA¹⁷ recommendations are also given in Table 3. They more or less represent averages of the given literature values, not including the most recent data by Seeley et al.¹¹ The averaged rate constant determined in the present work agrees within the error limits of the recommendations but our value is about 15% higher. This is expected to have a considerable influence on atmospheric chemistry models. In view of the more or less indirect procedure of deriving the rate constants adopted here, we will demonstrate that this difference cannot be explained by systematic deviations introduced in the data analysis.

For this purpose a nonlinear least-squares fit was performed with two of the decay curves from Figure 1 in which the rate parameters from eqs 8–10 are determined directly together with the total amplitude and the background signal (again omitting data points for $t \leq 0.2$ ms and after correction according to eq 5. From these parameters, which were found to be identical to those obtained using the alternative procedure, the corresponding model curves were constructed and subtracted from the corrected experimental curves. The residuals, normalized to a total amplitude of the decay curves of unity as in Figure 1, are displayed in Figure 4. The deviations are of the order of 0.2%, independent of time. Then the fit was repeated with the parameters $^1k_{1,\text{HO}_2}$ and $^1k_{\text{f,OH}}^1k_{\text{f,HO}_2}$ multiplied by a factor of 0.88 (the ratio 8.5/9.7 of the rate constants under consideration) and held fixed, while $^1k_{1,\text{OH}}$ (i.e., k_2), the total amplitude and the background were adjusted. Note that an OH yield of unity for the HO₂ + NO reaction is assumed in this procedure. The fitted OH loss rate coefficients showed only a small variation, resulting in changes in k_2 of -4% and -1.2% for curves (a) and (d), respectively. However, although the fitted decay curves appear quite similar to those obtained before at first glance, the curve parameters were found to have changed significantly by -4% and -5% (τ_1^{-1}), -12% and -9% (τ_2^{-1}), and +14% and +19% (c_1/c_2) for curves (a) and (d), respectively. Moreover, the quality of the fits decreased markedly, increasing χ^2 values by factors of 9.5 (a) and 8.3 (d). This is also reflected in the corresponding residuals displayed in Figure 4. It was checked that these deviations can not be attributed to the initial procedure of correcting the amplitudes.

Fixing only the value of $^1k_{1,\text{HO}_2}$ (i.e., k_1 to $8.5 \times 10^{-12} \text{ cm}^3 \text{ s}^{-1}$) resulted in a stronger change in k_2 of -30% and -18% for curves (a) and (d), respectively, while the OH yield of the HO₂ + NO reaction dropped to 87% in both cases. However,

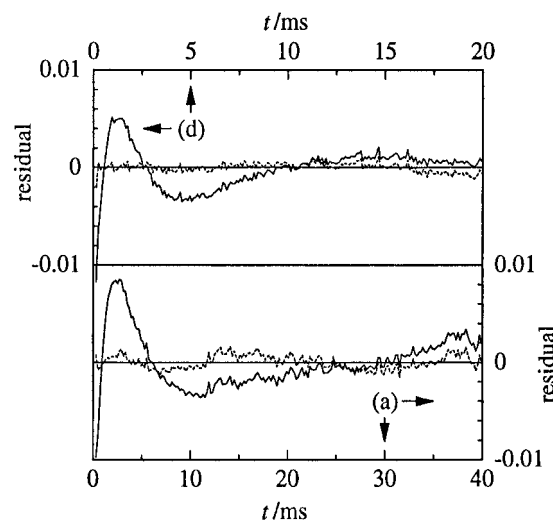


Figure 4. Residuals of nonlinear least square fits to decay curves (a) and (d) from Figure 1. Dotted line: result of a five-parameter fit including the total amplitude, the reaction rate parameters from eqs. 8–10 and the background signal. Full line: result of the same fit with the two parameters concerning the HO₂ + NO reaction scaled by a factor of 0.88 and held fixed (see text). Five-point averages of the data originally recorded are displayed.

the χ^2 values remain higher by factors of 5.4 (a) and 3.5 (d) than in the original fit.

From these results it is evident that the experimentally obtained curves are inconsistent with a rate constant significantly below $9.7 \times 10^{-12} \text{ cm}^3 \text{ s}^{-1}$ for the HO₂ + NO reaction. The reproducibility of this result is reflected in the scatter of the data in Figures 2 and 3 and in the statistical errors of the linear regressions given in Table 2. Moreover, the independence of the rate parameters concerning the HO₂ reaction on the H₂O₂ concentration supports the validity of the kinetic model employed. NO concentrations are estimated to be precise within 5%. Two different gas mixtures were used in combination with two different flow controllers without any noticeable effect. However, considering the suspected NO losses due to the chain reaction, background signal fluctuations, the stated error of the amplifier time constant, and uncertainties in the determination and constancy of τ_0^{-1} , we estimate a relative error of 10% in addition to the statistical error limits stated in Table 2. The quantum yield of OH formation in the 248 nm photolysis of H₂O₂ has not been considered here as a possible error source, as will be discussed below. At last our result is confirmed independently in the second part of the paper,³⁰ where a different experimental technique is applied and the rate constant has also been studied temperature dependent.

The OH yield of unity found for the HO₂ + NO reaction shows that neither channel 1b or 1c nor an isomerization of a possible HOONO intermediate to HNO₃ is of importance under the experimental conditions of this work. Since no deviation from a biexponential OH decay behavior has been observed, HOONO is either not formed or it dissociates to OH + NO₂ within less than ≈ 1 ms. The latter is possible in view of a recently calculated dissociation energy of only 66 kJ mol⁻¹ for the O–O bond in HOONO.³¹ However, since no indication for a pressure dependence of the rate constant was found, a complex mechanism for the reaction appears to be unlikely. The OH yields of the reaction were determined from correlations of the rate coefficients in Figure 2. They are therefore widely unaffected by the possible errors discussed for the rate constants, for example by the uncertainties in NO concentrations. Thus, a lower limit of 95% for the OH yield is derived from the stated statistical errors in Table 2.

The rate constant of the HO₂ self-reaction has been shown to depend on the presence of water vapor.¹³ This effect has been attributed to a complex formation followed by a displacement reaction.³² An HO₂-H₂O complex may also exhibit different reactive properties with respect to NO. However, under our experimental conditions we did not observe a significant effect of water vapor on the HO₂ loss rate coefficient. At the highest total pressure of 100 kPa, the water vapor mixing ratio of 0.017 used here corresponds to a partial pressure of 1.7 kPa, which is roughly half of the vapor pressure at room temperature (i.e., 50% relative humidity). The enhancement of the HO₂ + HO₂ reaction is about a factor of 2 under these conditions.¹³ We therefore conclude that the presence of water vapor has no influence on the HO₂ + NO reaction in the atmosphere.

The thorough work by Schiffman et al.²³ on the OH quantum yields of the 193 and 248 nm photolyses of H₂O₂ and HNO₃ has been carried out considering a large number of experimental parameters. Moreover, the vibrational Einstein transition probability used to calculate OH radical concentrations has been stated to be precise within 10%.²³ However, the reason for the measured low OH quantum yield²³ in the 248 nm photolysis remains unknown. The data obtained here confirm the more precise results on H-atom formation by Stief and DeCarlo²⁷ ($\phi_{\text{H}} + \phi_{\text{HO}_2} \leq 0.02$) and Vaghjiani and Ravishankara²⁸ ($\phi_{\text{H}} \leq 0.0001$). Energetically, formation of O(¹D) or O(³P) + H₂O and H₂ + O₂ is also possible.²³ However, these products are not expected to influence the present reaction system and their formation has been excluded experimentally.^{27,28}

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References and Notes

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